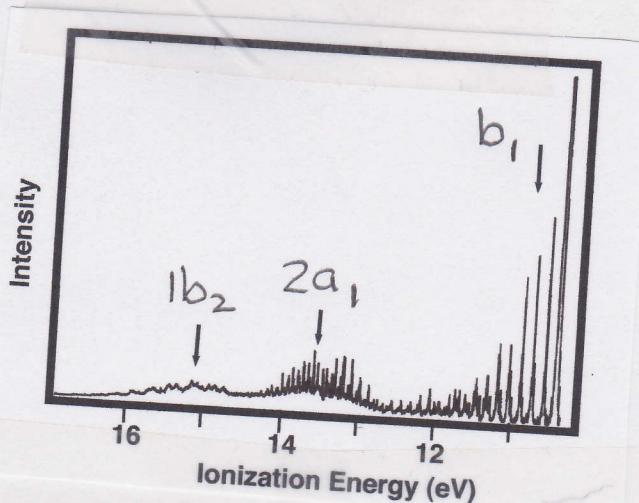
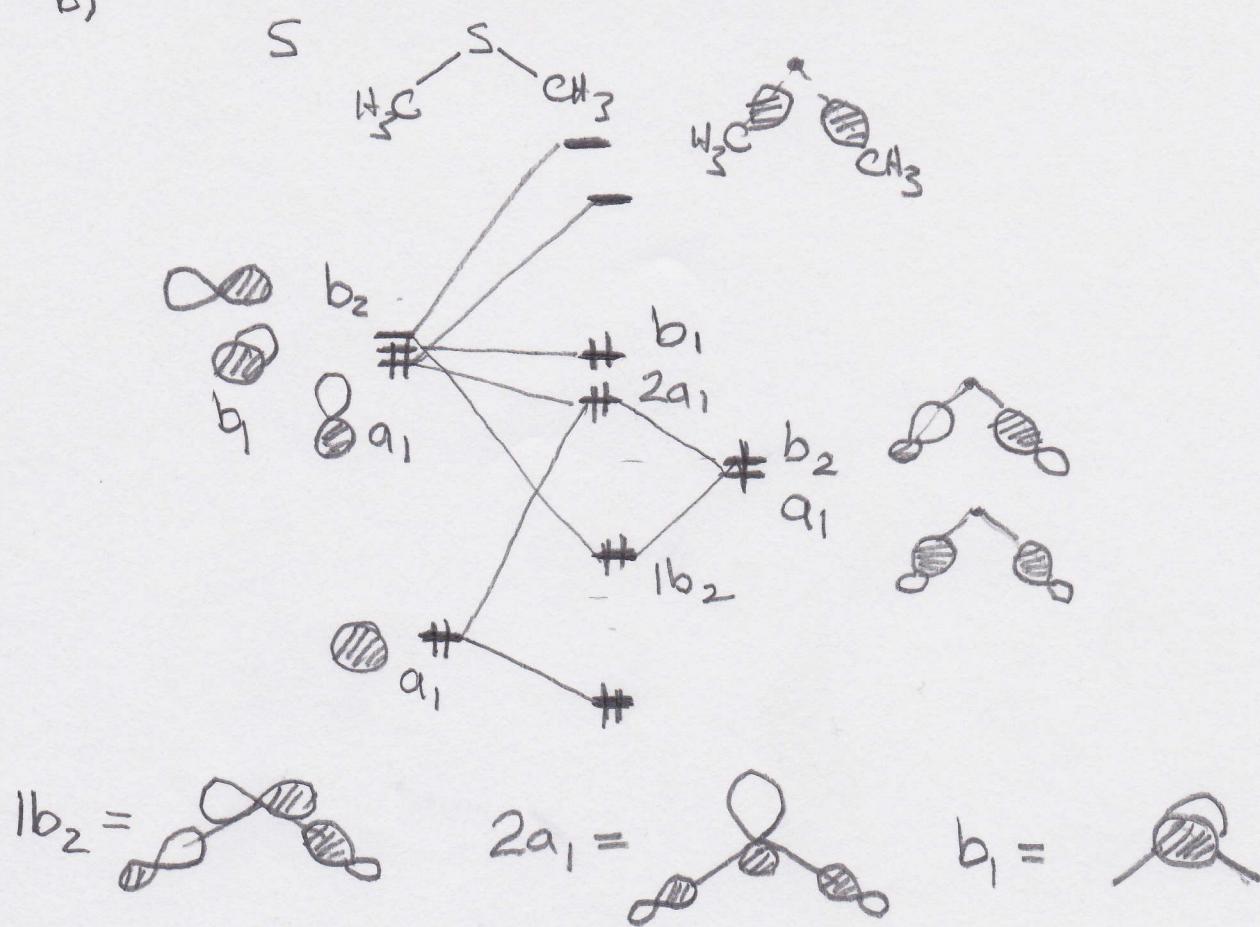


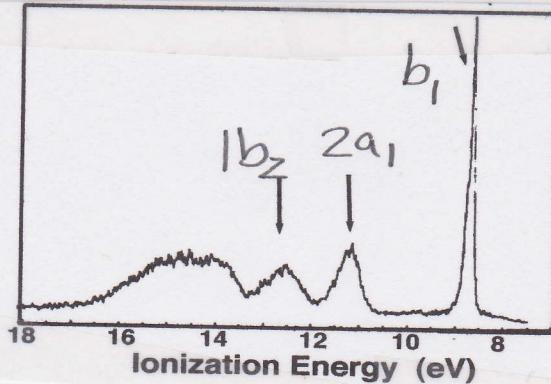
Answers - Chapter 7

1. a) the orbitals of H_2S should qualitatively be similar to those for H_2O - the PE spectrum is shown in Figure 7.9. Therefore,



b)

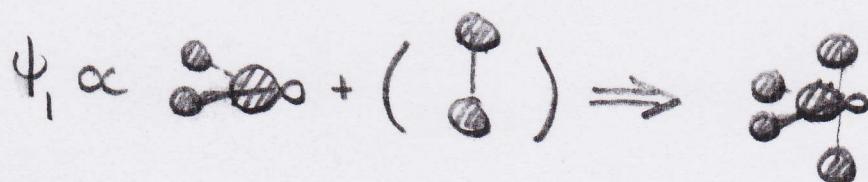
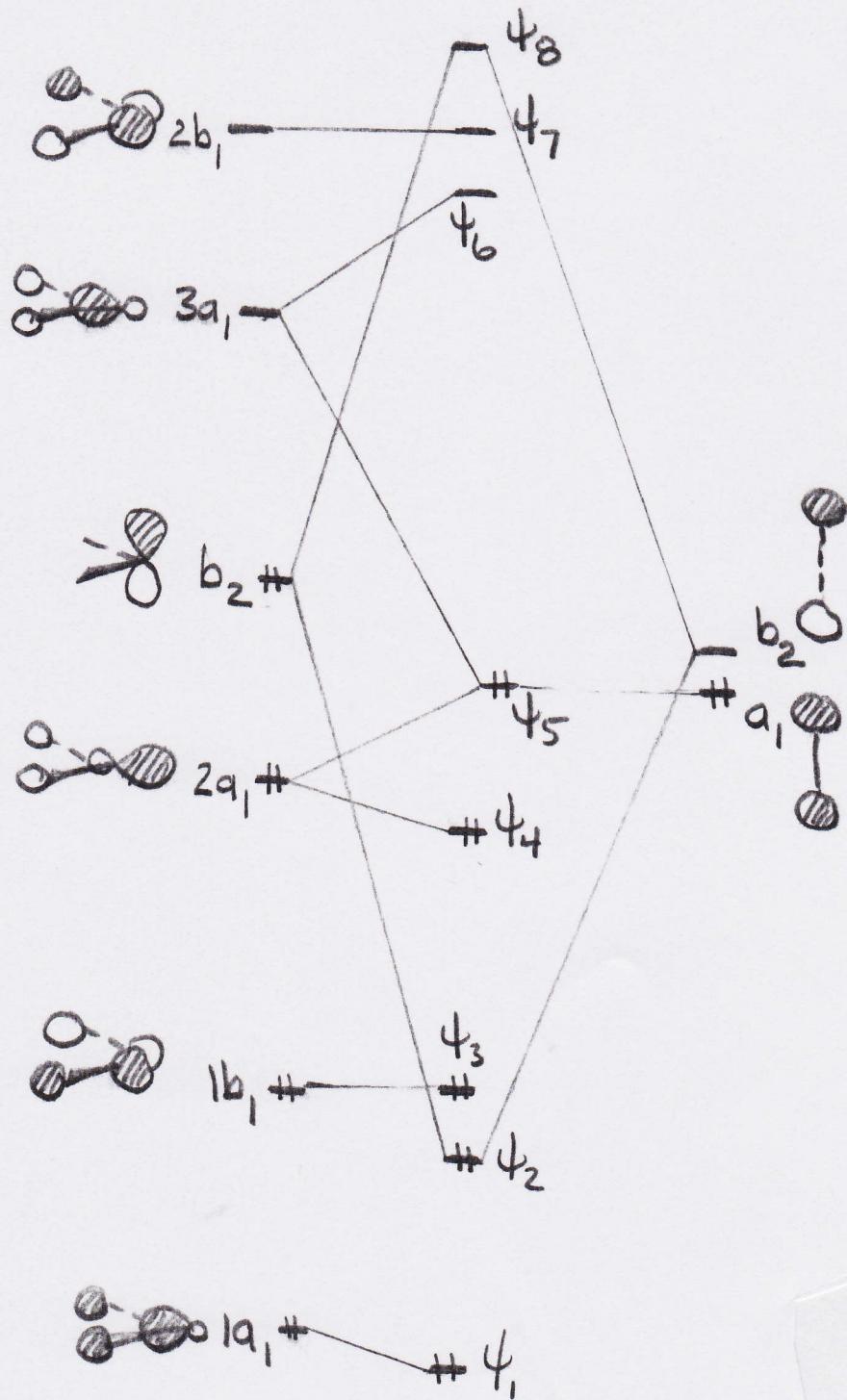
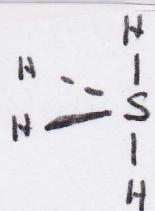
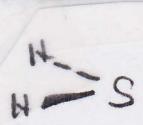




The first ionization is sharper (the vibrational progression is shorter) compared to the second ionization just as it is in H_2O and H_2S . This is consistent with b_1 MO being localized on S, as opposed to $2a_1$, which contains some CH_3 (or H) character and consequently the C-S-C (or H-S-H) bond angle changes in the radical ion state.

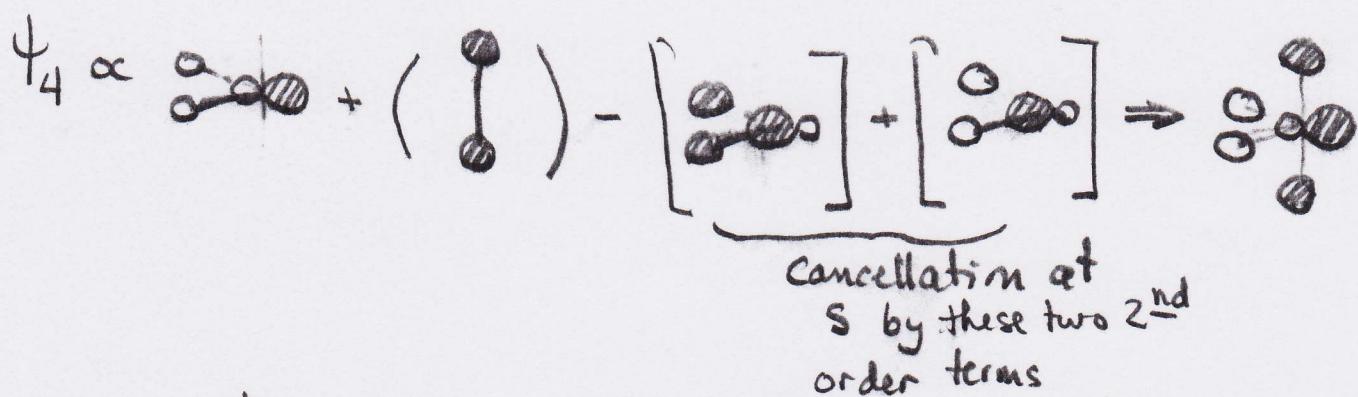
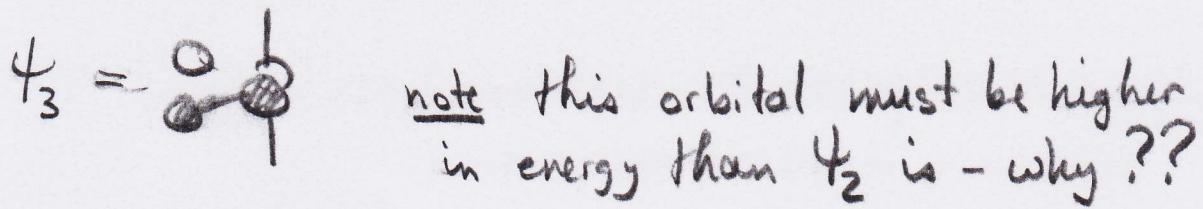
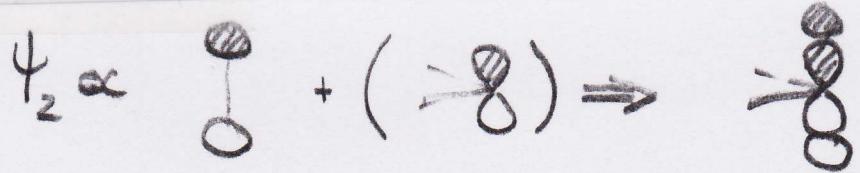
- c) The methyl group donates electron density compared to H. This pushes the MOs to higher energy or lower ionization potential in the PE spectrum.
- d) Going from $(\text{CH}_3)_2\text{S}$ to $[(\text{CH}_3)_3\text{C}]_2\text{S}$ the C-S-C bond angle is expected to open up. Recall from the Walsh diagram for H_2S in Figure 7.5; as the H-S-H angle increases, the $2a_1$ orbital is destabilized. Therefore, one would expect a lower ionization potential to be associated with it.
- The actual mechanism for the destabilization of $2a_1 \rightleftharpoons b_1$ is documented more fully in Chapter 10.5.

2.

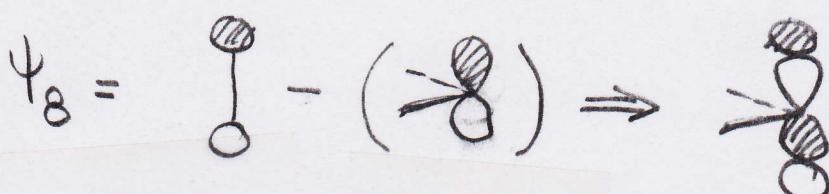
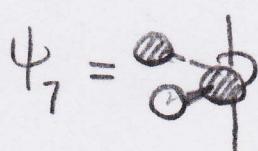
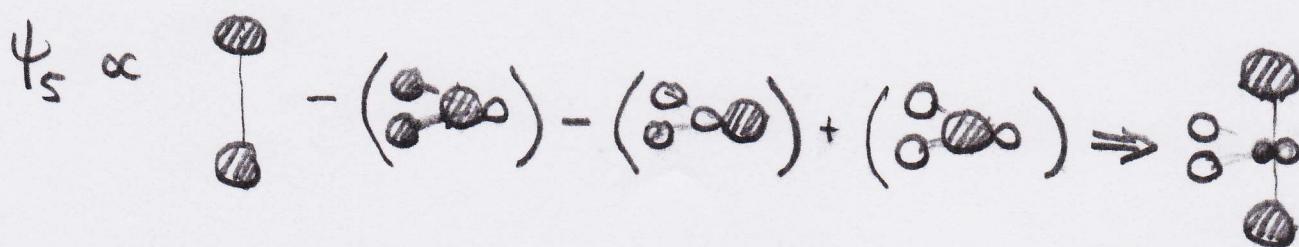


neglecting 2nd
order mixings with
2a₁ and 3a₁

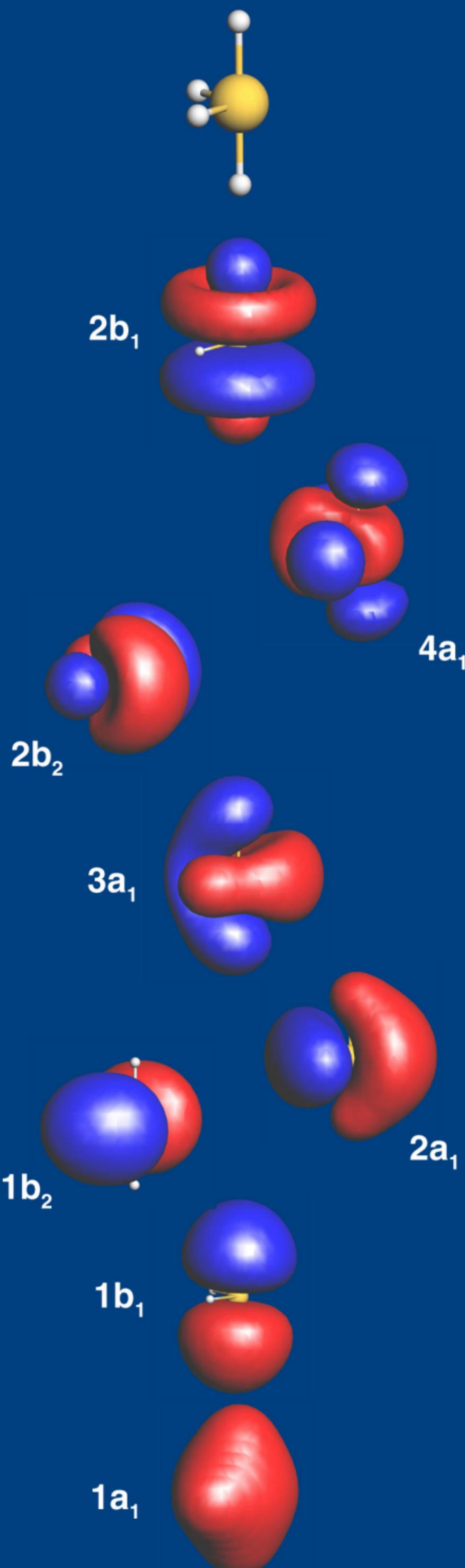
The MO's from an
Extended Hückel calculation are on the next page



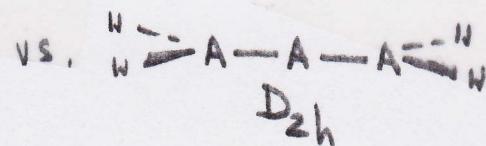
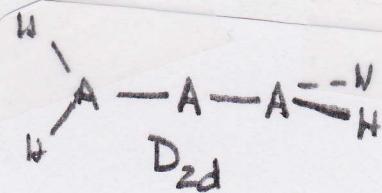
note overlap between $H_2 \sigma$ is determined by S A.O. at sulfur



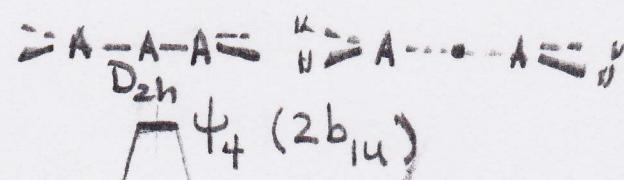
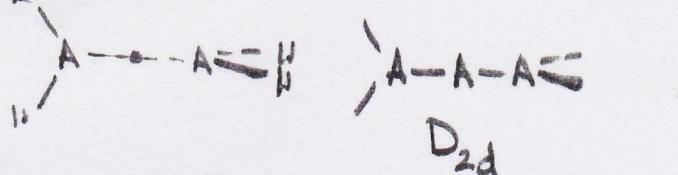
The MOs from an extended Hückel calculation are plotted on the next page.



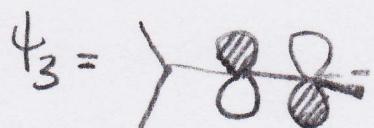
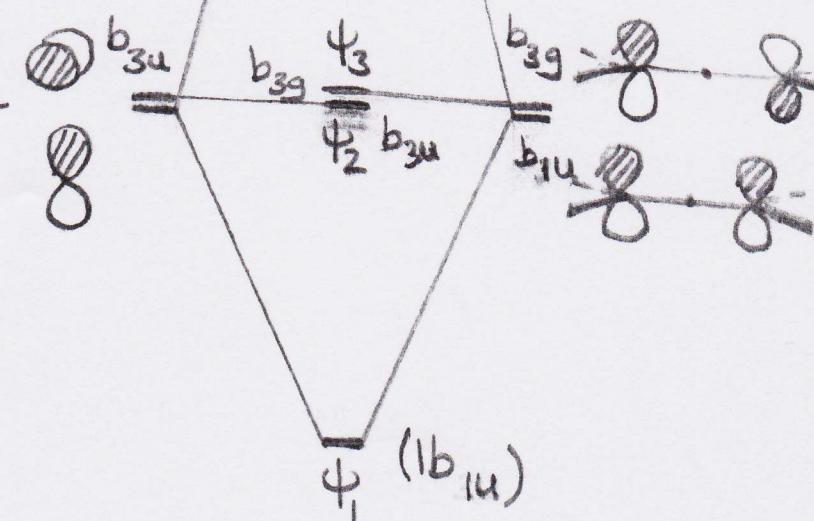
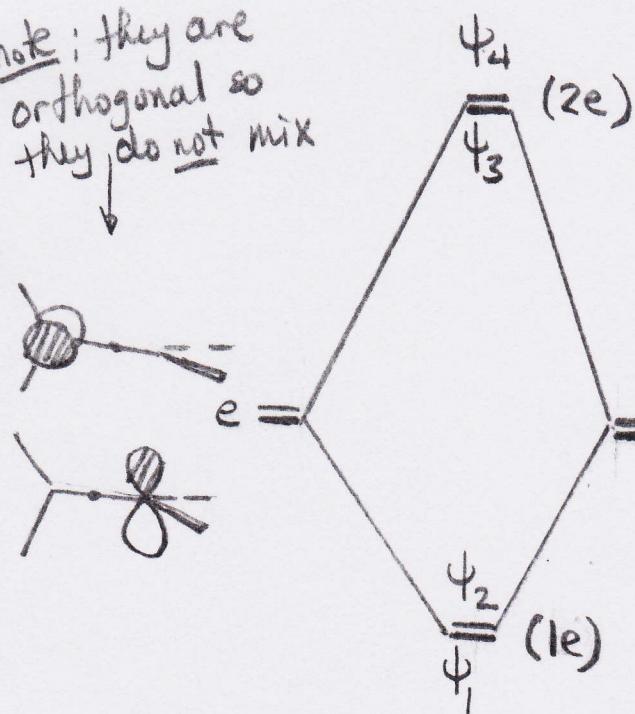
3.



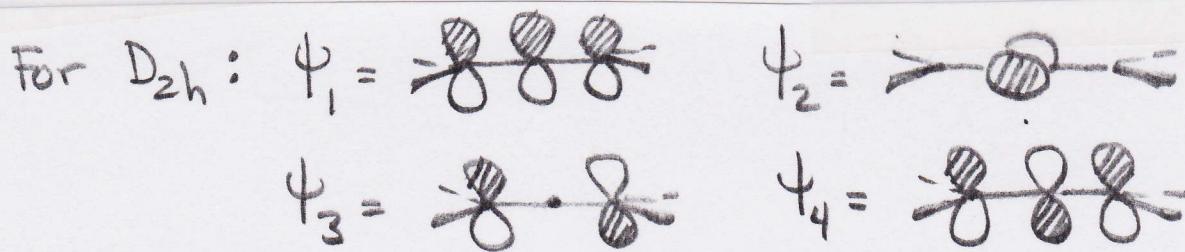
The easiest way to derive the π orbitals is via interacting $\text{H}_2\text{A} \cdots \text{AH}_2$ with the central A atom:



Note: they are orthogonal so they do not mix



} just the simple π/π^* levels!



Note $\psi_1, \psi_3 \in \psi_4$ are the π orbitals of an allyl system.

The π orbitals here are strictly analogous to the σ orbitals of H_2 and linear H_3 . Notice that

$1b_{1u}$ in D_{2h} is lower in energy than the $1e$ set in D_{2d}

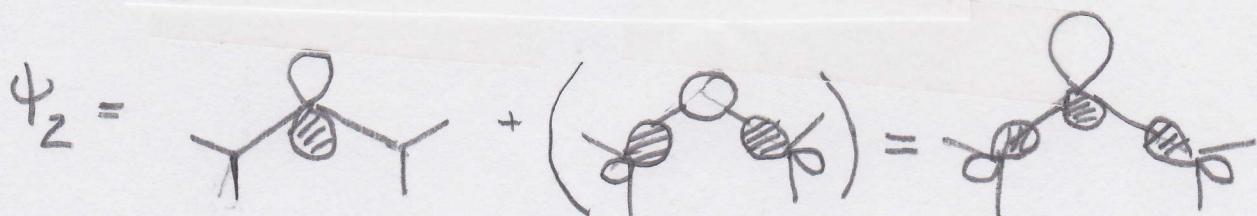
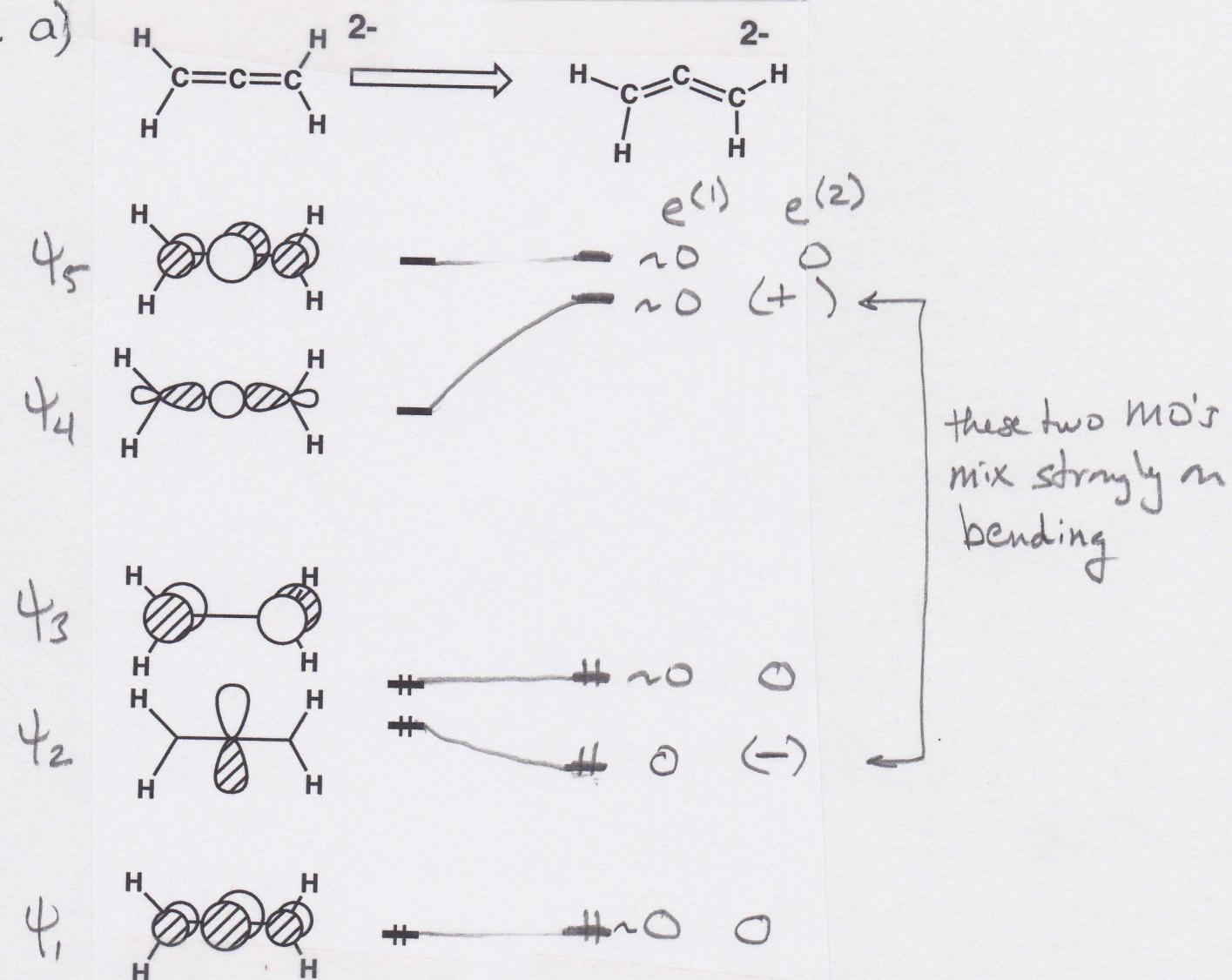
and $2b_{1u}$ is much higher in energy than the $2e$ set.

This is a straightforward application of Fig. 3.2. p. 56!

Therefore, the preferred structures are:

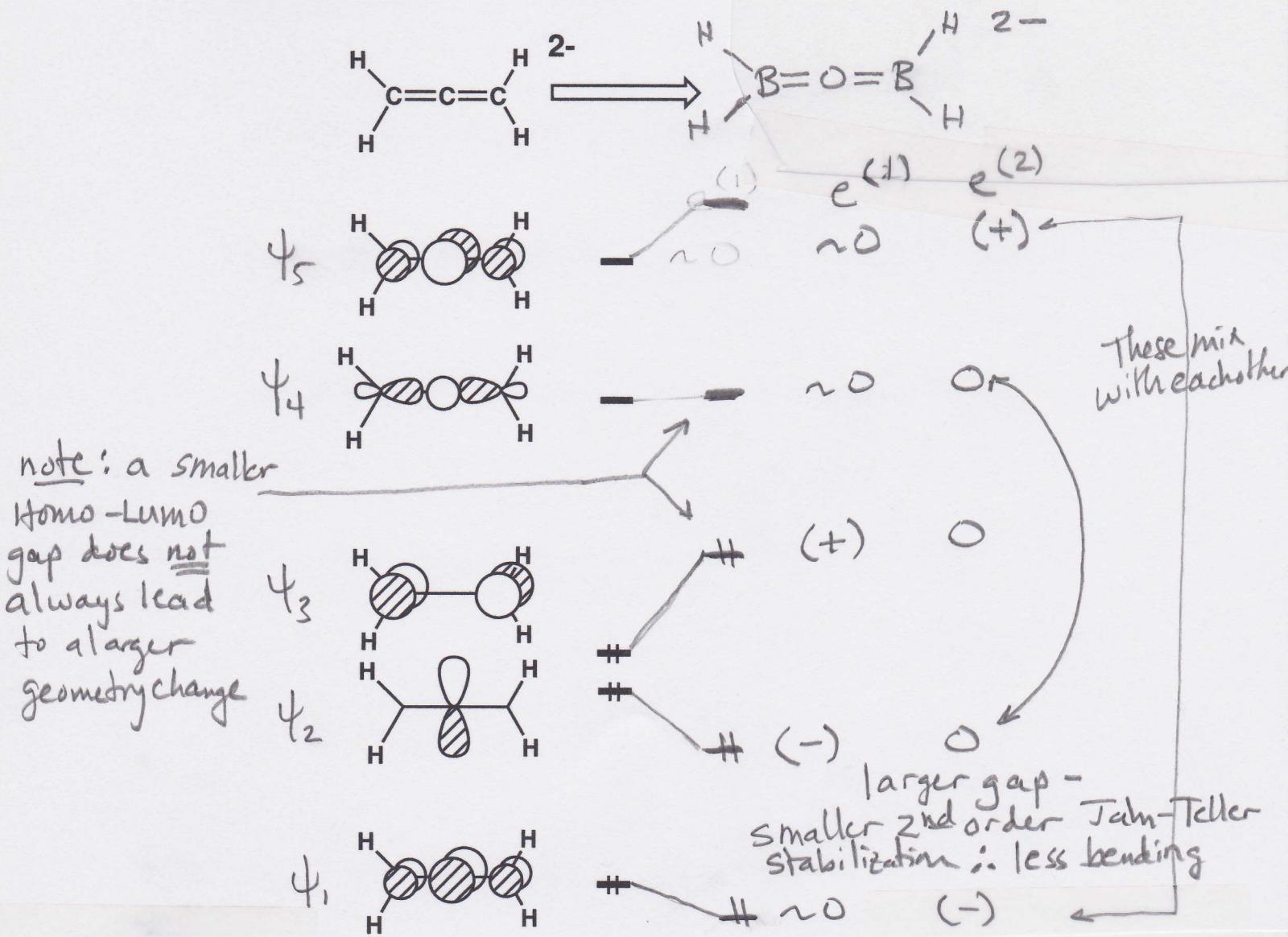
<u>number of e⁻</u>	<u>structure</u>
2	$D_{2h} \quad (1b_{1u})^2 \text{ vs. } (1e)^2$
4	$D_{2d} \quad (1b_{1u})^2 (b_{3u})^2 \text{ vs. } (1e)^4$
6	$D_{2h} \quad (1b_{1u})^2 (b_{3u})^2 (b_{3g})^2 \text{ vs. } (1e)^4 (2e)^2$
8	$D_{2d} \quad (1b_{1u})^2 (b_{3u})^2 (b_{3g})^2 (2b_{1u})^2 \text{ vs. } (1e)^4 (2e)^4$

4. a)



Note $\langle \psi_2 | \psi_4 \rangle = (+)$ as drawn

b) For B $\delta\alpha = (+)$; for O $\delta\alpha = (-)$. Either $\delta\alpha$ for two B atoms is greater in absolute magnitude than one O atom or they will cancel. Let's assume the latter. Then



5. See Figure 9.12 and the discussion around it for the answer. We shall see another H_2 complex in 14.47 and in 15.25.